An Analysis of Whistler Mode Radiation From the Spacelab 2 Electron Beam

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During the shuttle's Spacelab 2 mission the University of Iowa's plasma diagnostics package (PDP) was released from the shuttle to free fly. At times during this free flight when the PDP was magnetically connected to the shuttle, Stanford's fast-pulsed electron generator, located in the shuttle cargo bay, ejected a 1-keV, 50-mA electron beam. The plasma wave instrument on board the PDP detected intense whistler mode radiation during these beam ejections. This paper presents a study of a whistler mode emission detected during one particular continuous electron beam firing. Calculations indicate that the beam radiated approximately 1.6 mW in the whistler mode as the beam traversed the 200 m from the shuttle to the PDP. The emissivity also decreased by about a factor of 10 over this same distance. The measured wave powers are 10^7 greater than wave powers expected from incoherent Cerenkov radiation, verifying that the radiation is generated by a coherent process. Estimates of the emissivity based on measured electric field intensities in the beam indicate that the whistler mode noise is produced by radiation from electron bunches created by an electrostatic beam-plasma instability.

INTRODUCTION

In this paper we describe an electron beam experiment performed on the Spacelab 2 (SL 2) mission. The Spacelab 2 flight, which was launched on July 29, 1985, included an electron accelerator called the fast-pulsed electron generator (FPEG) from Stanford University and a spacecraft called the plasma diagnostics package (PDP) from the University of Iowa. During a 6-hour period on August 1, 1985, the PDP was released from the shuttle to investigate plasma effects in the vicinity of the shuttle. During the PDP free flight, the shuttle was maneuvered so that the PDP passed near magnetic field lines connected to the shuttle. Four such magnetic conjunctions were achieved. During one of these magnetic conjunctions a 1-keV, 50-mA electron beam was continuously ejected from the shuttle so that radiation effects could be monitored as the PDP passed near the magnetic field line carrying the beam. Plate 1 shows a frequency versus time spectrogram from the PDP plasma wave instrument during this electron beam event. The funnel-shaped signal extending from the electron cyclotron frequency fc down to approximately 30 kHz is whistler mode radiation from the electron beam. This whistler mode radiation was first described by Gurnett et al. [1986]. Our objectives in this paper will be to determine the total radiated power from the beam and compare it with the power predicted by various whistler mode radiation mechanisms.

Whistler mode emissions are known to be produced by both artificial and natural electron beams. Some of the early observations of emissions from artificial electron beams were made in the early 1970s by the University of Minnesota group using the Echo rocket experiments [Cartwright and Kellogg, 1974; Kellogg et al., 1976; Monson et al., 1976; Winckler, 1980]. Observations of beam-generated emissions were also made on the joint Franco-Soviet Artificial Radiation and Auroras Between Kerguelen and the Soviet Union (ARAKS) beam experiments [Dechambre et al., 1980]. The first space shuttle electron beam experiment was performed in March 1982 as part of the Space Transportation System 3 (STS 3) mission. On this flight the PDP was carried on the shuttle remote manipulator arm, and the FPEG was used to produce an artificial electron beam. During electron gun firings a possible whistler mode signal was detected from the electron beam [Shawhan et al., 1984]. In December 1983 the Phenomenon Induced by Charged Particles Beams (PICPAB) experiment was carried as part of the Spacelab 1 mission. The PICPAB experiment was specifically designed to investigate radiation from electron beams injected in the ionosphere. As in the other previous experiments, whistler mode signals were again detected when the electron gun was fired [Beghin et al., 1984]. Whistler mode radiation is also produced in the auroral zone in association with the field-aligned electron beams that are responsible for the aurora [Gurnett, 1966]. This radiation is usually called auroral hiss. Both upward and downward propagating auroral hiss has been observed [Mosier and Gurnett, 1969]. The downward propagating auroral hiss is associated with downward moving electron beams with characteristic energies of a few hundred electron volts [Gurnett, 1966; Hartz, 1969; Gurnett and Frank, 1972; Lasare and Hoffman, 1976]. The upward propagating auroral hiss often has a V-shaped spectrum called a "saucer" [Smith, 1969; Mosier and Gurnett, 1969; James, 1976] or a "funnel" [Gurnett et al., 1983].
Plate 1. A frequency versus time spectrogram from the PDP plasma wave instrument showing intense emissions during a dc electron gun firing. The funnel-shaped structure that extends from the electron cyclotron frequency \( f_c \) to about 30 kHz is whistler mode radiation from the beam.
The characteristic frequency-time shape of the saucer or funnel is a propagation effect that occurs for whistler mode waves propagating near the resonance cone. Gurnett et al. [1986], in their initial analysis of the Spacelab 2 results, showed that a similar wave propagation effect was observed for the whistler mode signals generated by the SL 2 electron beam.

**Electric Field Polarization**

To confirm that the radiation from the Spacelab 2 electron beam is propagating in the whistler mode, we compare the electric field polarization with the polarization expected for the whistler mode. The whistler mode has a polarization that depends on the wave frequency $f$, the wave normal angle $\theta$, the cyclotron frequency $f_c$, and the plasma frequency $f_p$. Using cold plasma theory [Stix, 1962], the electric field and index of refraction vectors can be calculated as a function of $f, \theta, f_c$, and $f_p$. For the whistler mode the index of refraction varies strongly as a function of the wave normal direction. This variation can be represented by an index of refraction surface, $n(\theta)$, which is a surface that is defined by the loci of index of refraction vectors at different wave normal angles. Figure 1 shows the index of refraction surface for the whistler mode. At a limiting wave normal angle, known as the resonance cone angle $\theta_{\text{res}}$, the index of refraction goes to infinity. The resonance cone angle is given by $\tan^2 \theta_{\text{res}} = -P/S$, where $P = 1 - f_p^2/f_c^2$ and $S = 1 - f_p^2/(f_c^2 - f_p^2)$. As the wave normal approaches the resonance cone, the electric field $E$ becomes linearly polarized with $E$ parallel to $n$. In this limit the electric field is quasi-electrostatic and the group velocity $v_g$ is perpendicular to $E$ and $n$ (see Figure 1).

In a previous paper [Gurnett et al., 1986] the funnel-shaped frequency versus time pattern of the radiation from the SL 2 electron beam was explained as a frequency dependent propagation effect for whistler mode emissions from the electron beam. Figure 2 is a sketch demonstrating how this effect works. Consider, first, a point source of whistler mode radiation, with the radiation propagating near the resonance cone. As the wave frequency increases, the resonance cone angle decreases, and the ray path direction $v_x$ becomes increasingly oblique to the magnetic field, approaching 90° as the frequency approaches the electron cyclotron frequency. Typical ray paths for a whistler mode emission propagating near the resonance cone are shown in the figure. As a spacecraft approaches the source, emissions near the gyrofrequency are detected first, since their ray paths are almost perpendicular to the beam. Progressively lower frequencies are then detected as the distance between the spacecraft and source decreases. This frequency dependent propagation effect creates a funnel-shaped emission pattern on a frequency versus time spectrogram. By comparing a modeled emission pattern using cold plasma theory to the funnel-shaped pattern shown in Plate 1, Gurnett et al. [1986] have confirmed that the whistler mode emission from the SL 2 electron beam is propagating upward from the shuttle near the resonance cone. The distinct funnel boundary is defined by the waves that propagate directly from the beam origin, the shuttle. The funnel on the spectrogram appears filled in, since the radiation is being emitted by a line source extending upward from the shuttle rather than a point source. Since the waves are upward propagating, $k \cdot v_x > 0$, where $v_x$ is the beam velocity vector.

In order to provide further confirmation that the radiation from the SL 2 electron beam is propagating near the resonance cone, an additional test was performed. This test compares model electric field directions in the PDP spin plane to their actual directions as measured by the PDP plasma wave instrument. To perform this test, a computer program was developed that calculates the angle $\phi$ between the projection of a model electric field onto the spin plane and a fixed reference direction. The fixed reference direction selected was the spin plane projection of the spacecraft-Sun vector. To compute $\phi$, the group velocity was assumed to be directed from a point on the beam toward the PDP with the electric field vector $E$ at an angle $\theta_{\text{res}}$ relative to the beam. This field geometry is the expected configuration for an upward propagating whistler mode wave near the resonance cone. Figure 3 shows the corresponding geometry of $E$, $v_x$, and $k$.

The electric field directions in the spin plane calculated using the model described above are compared to the measured electric field directions found from spin modulation maximums in the receiver data. The spin modulation maximums occur when the PDP electric antennas are aligned with the measured electric field in the spin plane, thus allowing a direct determination of this measured electric field direction.
the same direction as the beam \( \mathbf{k} \cdot \mathbf{v}_b > 0 \) provides a strong indication that the Landau resonance is involved. For example, the \( s = -1 \) cyclotron resonance produces radiation propagating in the opposite direction of the beam. Also, as will be discussed later, the Landau resonance gives the best agreement with the measured electric to magnetic field ratios.

To carry out the Poynting vector calculation we note that \( \mathbf{E} \) lies almost entirely in the plane defined by \( \mathbf{n} \) and the geomagnetic field (see Figure 3). The electrostatic and electromagnetic components of \( \mathbf{E} \) are given by \( E_x \cos \Delta \theta \) and \( E_x \sin \Delta \theta \), respectively, where \( \Delta \theta \) is the angle between \( \mathbf{E} \) and \( \mathbf{n} \) and \( E_x \) is the amplitude of the total electric field. The angle \( \Delta \theta \) is determined by the Landau resonance condition and cold plasma theory. The Landau resonance condition specifies the component of \( \mathbf{n} \) parallel to the geomagnetic field, i.e.,

\[
\mathbf{n}_\parallel = n \cos \theta = c/v_b \tag{1}
\]

where \( c \) is the speed of light. For a 1-keV electron beam moving parallel to the magnetic field, \( n_\parallel \) is approximately 15.9. A program was written that solves equation (1-20) of Stix [1962] for the magnitude and directions of \( \mathbf{n} \) and \( \mathbf{E} \). Equation (1-20) of Stix is

\[
\begin{bmatrix}
S - n^2 \cos^2 \theta & -iD & n^2 \cos \theta \sin \theta \\
iD & S - n^2 & 0 \\
n^2 \cos \theta \sin \theta & 0 & P - n^2 \sin^2 \theta
\end{bmatrix}
\begin{bmatrix}
E_x \\
E_y \\
E_z
\end{bmatrix} = 0
\]

where

\[
S = \frac{1}{2}(R + L) \tag{3}
\]

\[
D = \frac{1}{2}(R - L) \tag{4}
\]

\[
R = 1 - \sum_k \frac{\omega_{ek}^2}{\omega} \left( \frac{\omega}{\omega + \nu_b \omega_{ek}} \right) \tag{5}
\]

\[
L = 1 - \sum_k \frac{\omega_{ek}^2}{\omega} \left( \frac{\omega}{\omega - \nu_b \omega_{ek}} \right) \tag{6}
\]

Figure 4 shows the results of this comparison at four frequencies: 562, 311, 178, and 100 kHz. This figure shows the phase angle \( \phi \) between the projected electric field and the Sun vector as a function of time. The dots represent the modeled electric field directions computed assuming a resonance cone propagation scheme, while the crosses represent the measured electric field directions. The close agreement between the computed and measured electric field directions provides strong confirmation that the waves are propagating near the resonance cone and in the beam direction (i.e., \( \mathbf{k} \cdot \mathbf{v}_b > 0 \)), as indicated in Figure 3.

### Emitted Power

In this section we estimate the total power radiated from the beam in the whistler mode. By comparing the radiated power to the total power in the beam, the efficiency of the wave-beam interaction can be estimated and compared with various generation mechanisms. The power emitted from the beam in the whistler mode can be estimated by integrating the Poynting flux over a surface surrounding the beam. An inherent difficulty with this calculation is the determination of the phase and magnitude of the electric and magnetic fields in the Poynting flux, \( \mathbf{S} = \mathbf{E} \times \mathbf{H} \). Since three-axis measurements are not available and since phase measurements were not made, the Poynting vector cannot be determined directly. Computing the Poynting flux from the electric field alone is complicated by the fact that near the resonance cone the whistler mode is quasi-electrostatic and the ratio of the electromagnetic to electrostatic components is an exceedingly sensitive function of the wave normal angle. To compute the wave normal angle, we assume that the radiation is produced by the Landau resonance, i.e., \( \omega/k \parallel = \nu_b \). Since the beam velocity is known, this condition gives a well-defined value for the wave normal direction. The fact that the radiation is propagating in the same direction as the beam \( \mathbf{k} \cdot \mathbf{v}_b > 0 \) provides a strong indication that the Landau resonance is involved. For example, the \( s = -1 \) cyclotron resonance produces radiation propagating in the opposite direction of the beam. Also, as will be discussed later, the Landau resonance gives the best agreement with the measured electric to magnetic field ratios.

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The quantities $\omega_{pk}$ and $\omega_{ch}$ are the plasma and cyclotron frequencies for species $k$, while $s_k$ is the sign of the charge of species $k$. Using this program, we can calculate $n$ and $\Delta \theta$ at a particular wave frequency consistent with $n$. Since $\Delta \theta$ is now determined, the electrostatic and electromagnetic components of $E$ can be calculated. The calculated $\Delta \theta$ values are small, typically ranging from $0.06^\circ$ to $1.1^\circ$ from 31.1 to 562 kHz, and indicate that the wave is nearly electrostatic. It is easy to show that the magnitude of the Poynting vector is given by

$$P = 1 - \sum_k \frac{\omega_{pk}^2}{\omega^2}$$  \hspace{1cm} (7)

where $A = 1 - \cos^2 \Delta \theta$ and $B = \sin \Delta \theta \cos \Delta \theta$. In the derivation of (8), Faraday's law, $\mathbf{n} \times \mathbf{E} = c \mathbf{B}$, was used to eliminate the magnetic field in the $\mathbf{E} \times \mathbf{H}$ term. Consequently, $|S|$ is dependent only on $n$ and $E$, which are well-defined quantities. Note that as $\theta$ approaches the resonance cone angle, $n$ and $E$ become parallel and $|S|$ goes to zero. This behavior near $\theta_{res}$ is similar to an expression derived by Mosier and Gurnett [1971] in their paper addressing Poynting flux measurements of VLF hiss emissions.

Figure 5 shows the PDP trajectory relative to the beam. As can be seen, the PDP trajectory was nearly perpendicular to the beam. At closest approach, it passed within about 3 m of the beam at a distance of roughly 200 m from the shuttle. To compute the total radiated power, the Poynting flux is integrated over an imaginary surface perpendicular to the beam that includes the PDP trajectory. Assuming that the sampled intensities along this trajectory are constant around an annular ring of the area $dA$, centered on the beam, we can obtain the radiated power from the beam segment by evaluating the integral $P = \int S_y 2\pi R dR$, where $S_y$ is the field-aligned component of the Poynting vector and $R$ is the perpendicular distance from the beam to the PDP. Note that the evaluation of this integral will yield two values for the radiated power: one value from the inbound pass where the limits of integration extend from $R = -\infty$ to $R = 0$ and one value from the outbound pass where the limits of the integration now extend from $R = 0$ to $R = -\infty$. Figure 6 shows the average power spectral density from these two passes as a function of frequency. The error bars shown in the figure represent the standard deviations of the power values. Note that the power spectral density $dP/df$ is of the order of $10^{-9}$ W/Hz in the frequency range extending from 30 kHz to 1 MHz. The total emitted power in the 200-m beam segment from the shuttle to the PDP is found to be $P = 1.6$ mW. On the assumption that the power is emitted uniformly along the beam, the radiated power per unit length, $dP/dl$, is approximately $8 \times 10^{-6}$ W/m. Since the total power of the beam was 50 W, the beam converted approximately $3.2 \times 10^{-5}$ of its power to whistler mode radiation in the first 200 m. If the beam dissipated energy only via the whistler mode, the beam would only prop-
geometry, (9) can be rewritten as

$$\mathbf{n} \times \mathbf{E} = c \mathbf{B}$$

where \( \mathbf{E} \) is in the electric component and \( \mathbf{B} \) is in the magnetic component of the whistler mode waves. For the assumed field geometry, (9) can be rewritten as

$$nE_0 \sin \Delta \theta = cB_0$$

or

$$cB_0/E_0 = n \sin \Delta \theta$$

(10)

Using (10), we can compute \( n \sin \Delta \theta \) for various resonance conditions and compare this ratio with the measured \( cB/E \) ratio. The spectrum analyzer used with the PDP search coil can only provide measurements up to 178 kHz; therefore, the magnetic to electric field ratio can only be obtained in the 56-kHz, 100-kHz, and 178-kHz frequency channels. Also, the measured values of \( B \) at high frequencies using the search coil are highly uncertain because of inaccuracies in the calibration of the instrument. The preflight calibration was performed by placing a calibration coil in the search coil and surrounding the system in a \( \mu \) metal can. A problem arises at high frequencies (>10 kHz), where frequency dependent capacitances and inductances affect the current and the expected value of \( B \) from the calibration coils. Unfortunately, postflight calibrations under more ideal conditions (specifically, without the \( \mu \) metal can) have failed to reproduce the preflight calibrations. Our current best estimates are that \( B \) (and \( cB/E \)) are accurate only to within a factor of 2-4 at high frequencies. The range of measured \( cB/E \) values lies between 1.3 and 15.3. Assuming a Landau resonance, \( n \sin \Delta \theta \) is computed to be 0.54, 0.52, and 0.54 for 56, 100, and 178 kHz, respectively. Note that these values lie just outside the range of measured \( cB/E \) values and fall in the range when considering the factor of 2-4 uncertainty in the calibrations. For an \( s = +1 \) cyclotron resonance, however, \( n \sin \Delta \theta \) is computed to be between 0.05 and 0.08 for an \( s = -1 \) cyclotron resonance.

As mentioned earlier, the electric and magnetic field measurements also provide further evidence that the whistler mode waves are generated via a Landau resonance process. This argument involves a comparison of computed and measured \( cB/E \) ratios. Assuming a specific resonance condition and using the solution of equation (1-20) of Stix [1962], we can compute a unique value for \( n \) and \( cB_0 \) from the beam-generated whistler mode radiation detected by the PDP during the SL 2 mission. The total power from a coherent source goes as \( N^2 \), where \( N \) is the number of particles in coherence. Common coherent sources are plasma instabilities, lasers, and radio antennas. Incoherent mechanisms involve particles that are radiating independently. The power from the individual radiators must be added to get the total power emitted; thus the total power emitted is

$${\mathcal{P}}_{\text{tot}} = \sum_{i=1}^{N} P_i$$

(11)

where \( P_i \) is the power emitted by the \( i \)-th radiator. The total power emitted by the beam-generated whistler mode radiation is shown in Figure 7. Note that the emissivity drops by a factor of 10 from 100 to 200 m along the beam. This decrease in emissivity indicates that the efficiency of whistler mode generation decreases with increasing distance along the beam and that the generation mechanism is capable of dynamic changes in tens of meters. If the emissivity continues to drop at the rate observed between 100 to 200 m, the radiation would be undetectable by the PDP at source distances more than about 1 km from the shuttle. This result may explain why Dynamics Explorer 1 (DE-1), which was magnetically connected to the shuttle during a gun firing on the STS 3 mission, did not see beam-generated whistler mode radiation in the vicinity of the streaming electrons [Inan et al., 1984]. From the SL 2 measurements it appears that strong whistler mode emissions are probably generated only in close proximity to the source of the beam.
Coherent Generation Mechanisms

As concluded in the previous section, some coherent process must be involved in the whistler mode wave generation from the SL 2 electron beam. Coherent processes can be divided into two classes: direct and indirect. Direct mechanisms involve the direct conversion of energy from an unstable particle distribution to electromagnetic radiation, whereas indirect mechanisms involve the intermediate generation of one or more electrostatic modes which are coupled to the escaping electromagnetic radiation. This section will discuss possible direct and indirect mechanisms that may explain the whistler mode radiation.

Since an unstable electron distribution is present in the beam, the escaping electromagnetic radiation may result from direct conversion of the beam energy to electromagnetic radiations. Such a mechanism has been proposed by Maggs [1976] for the generation of auroral hiss. In his model, incoherent Cerenkov radiation produced by an auroral electron beam is directly amplified via a whistler mode plasma instability within the beam. It seems reasonable that this wave generation mechanism could be applied to the whistler mode waves emitted from the SL 2 electron beam; however, a problem arises in doing so. Unlike auroral beams the path length for wave growth in the SL 2 beam is very short, only 2–3 electron cyclotron radii (6–9 m). Using the Landau resonance condition and the fact that the emission is propagating near the resonance cone, we can derive an expression for the wavelength of the whistler mode radiation,

\[ \lambda \approx (V_b/\omega_c) \cos \theta_{\text{esc}} \]  

(11)

which for the nominal parameters has a value of about 20 m. This wavelength is greater than the path length, which completely invalidates any mechanism involving exponential growth. Even if that were not the case, for typical whistler mode group velocities of \(10^7\) m/s the amount of time the wave spends in the beam is so short, only about \(10^{-6}\) s, that unreasonably high growth rates \(\omega > \omega_c \approx 10^6 \text{s}^{-1}\) would be required to generate the radiation. No whistler mode instability is known that can produce such large growth rates from realistic electron distribution functions. These same conclusions were also reached by Jones and Kellogg [1973] in their paper addressing the growth rates of whistler mode radiation from artificially created electron beams.

Next we consider mechanisms involving the intermediate generation of electrostatic waves in the beam. Any density perturbation or bunch created by an electrostatic wave in the beam is capable of emitting coherent Cerenkov radiation. The radiated power from this process will have a frequency spectrum similar to that of a single radiating electron; however,
the wave power will be much greater, since the emitted power goes as $N^2$, where $N$ is the number of electrons in a bunch. Coherent Cerenkov radiation from a bunched beam has been considered previously by Bell [1968].

Beam-plasma instabilities are known to be capable of creating intense electrostatic waves and density perturbations in the beam. We estimate the number of coherently bunched electrons required to account for the observed whistler mode radiation. The total power emitted from the electron bunches in the beam is $(dP/df)_\text{total} = (dP/df)_\Delta N \Delta \nu$, where $(dP/df)_\gamma$ is the power radiated by each electron, $\Delta N$ is the typical number of electrons in a bunch, and $\Delta \nu$ is the number of bunches in the 200-m segment of the beam. Beam-plasma instabilities are known to create an electrostatic wave near the local electron plasma frequency (3 MHz). Such an emission is, in fact, observed near 3 MHz [see Gurnett et al., 1986]. The corresponding wavelength is $\lambda_c = V_c/\Delta \nu \approx 7 \text{ m}$, which is assumed to be the approximate length of each bunch. This wavelength can then be used to calculate $\Delta \nu$, the number of bunches in the first 200 m of the beam. This number is $\Delta \nu \approx 29$. Once the radiated power from the 200-m beam segment, $(dP/df)_\text{total}$, is about $10^{-6}$ W/Hz, $(dP/df)_\gamma$ is about $10^{-29}$ W/Hz, and $\Delta \nu \approx 29$, each bunch must contain about $\Delta N = 2 \times 10^9$ electrons in order to account for the observed radiated power.

Next we estimate the required electric field strength for the electrostatic wave in the beam that forms the bunches. Assuming that the beam diameter is about 2 cyclotron radii, we can estimate the electron number density in the bunch using the formula:

$$\Delta n = \frac{\Delta N}{\pi r_c^2} \Delta L \quad (12)$$

where $\Delta L$ is the bunch length and $r_c$ is the cyclotron radius (2 to 3 m). The required number density is found to be about $\Delta n = 1 \times 10^7$ electrons/m$^3$. Note that the fractional density perturbation in the beam $\Delta n/n_0$ is only about 0.01; thus a relatively small perturbation can account for the measured whistler mode power. Poisson’s equation can be used to determine the magnitude of the self-consistent electric field needed to generate this density perturbation:

$$\Delta E/\Delta L = e\Delta n/r_c \quad (13)$$

From (13) we estimate that an electric field of the order of 1–2 V/m is needed to create the required coherence in the beam electrons.

Although the PDP did not fly directly through the beam during free flight, when the PDP was on the remote manipulator arm, it did provide measurements in the beam. During these times, intense field-aligned electric field signals near $f_p$ were measured with amplitudes greater than 0.3 V/m, sufficiently large to saturate the receiver. This value is within a factor of 10 of the required amplitudes needed for radiative coherence of the beam electrons. The good agreement between the calculated and measured electrostatic field strengths strongly suggests that electron bunches generated by a beam-plasma instability can account for the observed whistler mode power.

CONCLUSIONS

Like auroral hiss the whistler mode radiation from the SL 2 electron beam was found to be propagating near the resonance cone in the same direction as the beam. This conclusion was reached by comparing the measured and modeled electric field polarizations in the PDP spin plane. The total whistler mode power radiated by the beam between the shuttle and the PDP was estimated to be 1.6 mW. This represents a conversion of about $3.2 \times 10^{-4}$ of the beam power to whistler mode radiation in a path length of 200 m. The linear emissivity was also observed to decrease with increasing distance along the beam, thereby indicating that most of the whistler mode radiation is generated relatively close to the source of the beam.

The mechanism that best explains the radiated power and frequency range of the whistler mode radiation is Cerenkov radiation from coherently radiating electron bunches in the beam. The bunches are most likely created by an electrostatic instability in the beam. The estimated electric field amplitudes required to generate the bunching are consistent with the field strength of electrostatic noise measured in the beam.

Future work includes calculating the radiated power from a model electron beam obtained from a particle simulation. The calculated power can then be compared to the measured whistler mode power to confirm that coherent radiation from electron bunches generates the signal.

Acknowledgments. We are grateful to Bill Kurth and John Steinberg for their helpful comments concerning this work. We also like to thank Terry Averkamp for his valuable input on data analysis and John Birkbeck for a fine drafting job. Part of this research was funded by NASA Graduate Student Researchers Training grant NGT-50034. The research at the University of Iowa was also supported by NASA through contract NAS8-32807 and grants NGL 16-001-002 and NGL 16-001-043 and by the Office of Naval Research through contract N00014-85-K-0404. The research at Stanford University was supported by NASA through grant NAGW-235.

The Editor thanks R. H. Holzworth and P. L. Pritchett for their assistance in evaluating this paper.

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(Received February 27, 1987; revised July 23, 1987; accepted September 2, 1987.)