Interplanetary Particles and Fields, November 22 to December 6, 1977: Helios, Voyager, and Imp Observations Between 0.6 and 1.6 AU

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In the period November 22 to December 6, 1977, three types of interplanetary flows were observed: a corotating stream, a flare-associated shock wave, and a shock wave driven by ejecta. Helios 2, Imp 7, 8, and Voyager 1, 2 were nearly radially aligned at ≈ 0.6, 1, and 1.6 AU, respectively, while Helios 1 was at ≈ 0.6 AU and 35° east of Helios 2. The instruments on these spacecraft provided an exceptionally complete description of the particles and fields associated with the three flows and corresponding solar events. Analysis of these data revealed the following results. (1) A corotating stream associated with a coronal hole was observed at 0.6 and 1 AU, but not at 1.6 AU. The stream interface corotated and persisted with little change in structure even though the stream disappeared. A forward shock was observed ahead of the interface and moved from Helios 2 at 0.6 AU to Voyager 1, 2 and 1.6 AU; although the shock was ahead of a corotating stream and interface, the shock was not corotating, because it was not seen at Helios 1, probably because the corotating stream was not stationary. (2) An exceptionally intense type III burst was observed in association with a 2B flare of November 22. The exciter of this burst (a beam of energetic electrons) and plasma oscillations (presumably caused by the electron beam) were observed by Helios 2. (3) A nonspherical shock was observed in association with the November 22 flare. This shock interacted with another shock between 0.6 and 1 AU, and they coalesced to form a single shock that was identified at 1 and at 1.6 AU. (4) A shock driven by ejecta was studied. In the ejecta the density and temperature were unusually low, and the magnetic field intensity was relatively high. This region was preceded by a directional discontinuity at which the magnetic field intensity dropped appreciably. The shock appeared to move globally at a uniform speed, but locally, there were fluctuations in speed and direction of up to 100 km/s and 40°, respectively. (5) Three types of electrostatic waves were observed at the shocks, in different combinations. The detailed wave profiles differed greatly among the shocks, even for spacecraft separations of ≈ 0.2 AU, indicating a strong dependence on local conditions. However, the same types of fluctuations were observed at 0.6 and at 1.6 AU. (6) Energetic (50–200 keV) protons were accelerated by the shocks. The intensities and durations of the fluxes varied by a factor of 12 over longitudinal distances of ≈ 0.2 AU. The intensities were higher, and the durations were lower, at 1.6 than at 0.6 AU, suggesting a cumulative effect. (7) Energetic (≈ 50 keV) protons from the November 22 flare were observed by all the spacecraft. During the decay, Helios 1 observed no change in intensity when the interface moved past the spacecraft, indicating that particles were injected and moved uniformly on both sides of the interface. Helios 2 observed an increase in flux not seen by Helios 1, reaching maximum at the time that a shock arrived at Helios 2. The intensity dropped abruptly when the interface moved past Helios 2, indicating that the ‘extra’ particles seen by Helios 2 did not penetrate the interface.

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I. INTRODUCTION

It is customary to speak of three types of flows in the solar wind, which Chapman [1964] called streams, flare shells, and solar wind. We shall refer to these as (1) corotating streams (also known as corpuscular streams, stationary streams, and high-speed streams), (2) ejecta (also called nascent streams, flare ejecta, jets, active wind, drivers, and pistons), and (3) slow flows (also called quiet wind, ambient wind, and structureless wind), respectively. The term piston has been applied both to corotating streams [e.g., Dryer and Steinolfson, 1976] and to ejecta [e.g., Dryer et al., 1972]. Corotating flows are separated from slow flows by a thin boundary called a stream interface [Belcher and Davis, 1971; Burlaga, 1974], which in some cases is a tangential discontinuity [Burlaga, 1974; Gosling et al., 1978]. Ejecta are presumed to be separated from slow flows by a thin boundary called a contact surface [Lee and Chen, 1968; Dryer, 1980]; this too is called a piston by some authors [e.g., Dryer, 1980]. Both corotating streams and ejecta may be preceded by a shock.

There are numerous studies of the above phenomena based on observations from just one spacecraft, but these cannot separate spatial variations from temporal changes. Data from two or more widely separated spacecraft are needed to study nonstationary corotating streams, transient ejecta, and interacting flows. There have been relatively few such observational studies [e.g., Dryer et al., 1972; Gosling et al., 1976; Intriligator, 1976; Lazarus et al., 1979; Schwenn et al., 1980; Smith and Wolfe, 1977, 1979; Vaisberg and Zastenker, 1976]. Some attempts have been made to model multipoint observations [Gosling et al., 1976; Dryer et al., 1978a, b].

In the period November 22 to December 6, 1977, Helios 1 and 2, Imp 7 and 8, and Voyager 1 and 2 were aligned very favorably for the investigation of solar outputs (Figure 1), and during this period, which was part of Study of Travelling Interplanetary Phenomena (STIP) interval IV from October 15
to December 15, 1977, several significant solar events occurred. Recognizing that this interval (and a similar interval in September-October 1977) offered a unique opportunity for a comprehensive study of interplanetary shocks, flows, magnetic fields, and energetic particle phenomena, a workshop was organized to bring together experimenters from the Helios, Voyager, and Imp programs. The meeting was organized by S. M. Krimigis with the support of the Voyager and Helios team leaders. This paper is based on some of the results of that workshop. The purpose of this paper is to present a description and an analysis of the principal interplanetary events that were observed in the period November 22 to December 6, 1977, by Helios 1, 2, Voyager 1, 2, and Imp 7, 8.

Three flow systems were observed in the period under consideration: (1) a corotating stream and a stream interface associated with a coronal hole, (2) a shock wave and an energetic particle event associated with a 2B flare, and (3) an isolated shock wave whose origin is uncertain.

This paper is based on data from 28 experiments from six spacecraft. The experiments and the corresponding principal investigators are listed in Table 1. Nearly complete measurements of solar wind plasma, magnetic fields, and plasma waves are available from all spacecraft. Radio waves, plasma waves, and energetic electrons associated with the November 22 event are available from Helios 1, 2 and Voyager 1, 2. Data describing low-energy protons associated with the November 22 event are available from Helios 1, 2 and Voyager 1, 2.

We begin in section 2 by discussing the corotating stream and its associated shock and interface; this flow system was relatively simple, and the other two events interacted with it. Section 3 discusses the particles, fields, and flows associated with the flare of November 22. Section 4 analyses a relatively simple, isolated shock wave that passed all of the spacecraft in the early days of December 1977. Plasma waves at the shocks in the three events are discussed qualitatively in section 5. Energetic protons accelerated by the shocks and injected by the November 22 flare are described in section 6. Section 7 summarizes the results.

2. COROTATING STREAM, INTERFACE, AND SHOCK

A stream that was observed successively by Helios 1, Helios 2, Imp 7, 8, Voyager 1, and Voyager 2 is shown in Figure 2, which shows bulk speeds from the experiment of Rosenbauer on Helios 1, 2 and from the experiments of Bridge on Imp 7, 8 and Voyager 1, 2. Sixteen-minute averages of \( V \) are plotted versus time, and the phase is chosen such that the arrival time of the stream interface at each spacecraft is coincident with the vertical line marked 'interface.' The stream interface is readily identified as an abrupt decrease in density and an abrupt increase in temperature at the front of a stream [Belcher and Davis, 1971; Burlaga, 1974, 1975]. In this case the interface at each spacecraft can be seen in Figure 3, where the time profiles of 16-min averages of the density \( n \) and temperature \( T \) are plotted. Figure 2 shows that the interface and stream arrived at Helios 1 on November 23, at Helios 2 on November 25, at earth on November 27, and at Voyager 1 and 2 on November 29. The 2-day interval between successive encounters of the interface is approximately that which is expected for a 'corotating spiral' corresponding to a streamline with a speed of 400 km/s, as illustrated at the bottom of Figure 2.

The precise corotation times of the interface from one spacecraft to the next are shown in Table 2, together with the 'predicted' corotation times computed from the equation \( t_2 - t_1 = (r_2 - r_1)/V + (\phi_2 - \phi_1)/\Omega_s \), with allowance for the spacecraft motions (here \( \Omega_s \) is the sidereal rotation period of the sun; \( V \) is the solar wind speed; \( \phi_1 \) and \( \phi_2 \) are the longitudes of the spacecraft at time \( t_1 \) (when the interface passed the first spacecraft) and a later time \( t_2 \) (when the interface passed the second spacecraft), respectively; and \( r_1 \) and \( r_2 \) are the radial distances from the sun of the two spacecraft at \( t_1 \) and \( t_2 \).) Table 2 shows that the predicted corotation times are close to the observed corotation times, the difference being \( \leq 15\% \) in the three largest time intervals. These small differences may be due to small irregularities in the shape of the surface of the interface. Thus we conclude that the interface was a corotating feature, and we infer that the stream which followed it was likewise corotating.

The low densities in the stream (Figure 3) and the fact that it was corotating suggest that its source was a coronal hole [Hundhausen, 1977; Burlaga, 1979]. A coronal hole, tentatively identified in the Kitt Peak He 10830-Å maps, passed central meridian on November 24, 25. The observed peak speed of the stream in question was \( \approx 500 \) km/s; thus if its source was the coronal hole and if it propagated at nearly constant speed, the stream should have arrived at the earth on November 27, which, in fact, it did.

The dynamical evolution of the corotating stream in Figure

<table>
<thead>
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<th>Table 1: Principal Investigators</th>
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<tr>
<td>Plasma Analyzer</td>
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<tr>
<td>Helios 1</td>
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<td>Helios 2</td>
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<td>Voyager 1</td>
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<td>Voyager 2</td>
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<td>Imp 7</td>
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<td>Imp 8</td>
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rated the stream. The wave moving away from the sun (i.e., ahead of the stream) evolved into a forward shock (see discussion below). The importance of momentum flux in corotating stream dynamics has been discussed quantitatively by Pizzo [1980a, b] for some conventional stream profiles. A decrease of $V$ with increasing distance has been discussed by Dryer and Steinolfson [1976].

The structure of the stream interface observed by Voyager 1 and 2 is shown in Figure 4. In both cases the density and temperature transitions occurred in $\approx 30$ min, consistent with the durations of some of the interfaces observed at 1 AU by Burlaga [1974]. The $n$, $T$ profiles observed at Helios 1, 2 are very similar to those in Figure 4. The magnetic field intensity reached a maximum at the interface (see Figure 4), as is usually the case [Burlaga, 1974; Siscoe, 1972]. In this case there was a large change in magnetic field direction across the interface at both Voyager 1 and Voyager 2. It is significant that all of the parameters just described ($n$, $T$, $V$, and $B$) had nearly the same profile at Voyager 2 as at Voyager 1, despite the separation of $\approx 0.2$ AU; this shows that the internal structure of a stream interface can be coherent over a relatively large distance. Plasma wave observations at the interface at Voyager 2 (Figure 4) show no significant wave emission in the frequency range 10–562 Hz, suggesting that the interface was relatively stable. Similar observations of a different interface described by Garnett et al. [1979b] showed the same result.

A 'corotating shock' (which we label shock B) was observed by Voyager 1 and 2; this is shown at high resolution in Figure 5. The identification of the disturbance as a shock is based on the simultaneous, abrupt increases in $V$, $N_p$, $T_p$, and $F = |B|$ and on the simultaneous change in the characteristics of the plasma waves. The observation of a shock at Voyager 1 and 2 is not surprising, since models of corotating streams [e.g., Hundhausen, 1973; Hundhausen and Burlaga, 1975; Gosling et al., 1976; Steinolfson et al., 1975; Dryer et al., 1978b] predict the development of corotating shocks as streams evolve with distance from the sun, and many such shocks have been observed beyond 1 AU [Smith and Wolfe, 1977]. The shock normals computed from the Voyager plasma and magnetic field data using the method of Lepping and Argentiero [1971] were directed 9° and 14° west of radial, respectively (see Table 3 and Figure 2), consistent with corotation. At Voyager 2 the angle between the shock normal and the upstream magnetic field was 14.6°; the corresponding angle at Voyager 1 was 15.8°. The local shock speed was 400 ± 10 km/s relative to a fixed frame shock at both Voyager 1 and Voyager 2 (Table 3). This speed and the computed shock normals imply a time delay between Voyager 1 and Voyager 2 of 4.5 ± 1.3 hours. This compares favorably with the observed time delay of 5 hours 17 min.

Shock B probably passed Helios 2 and Imp 8 on November 25 and 26, respectively (see Table 3). This is significant, because shocks are rarely observed ahead of corotating streams at $\approx 1$ AU [Ogilvie, 1972]. The identification of shock B at Helios 1 is based on the observations that (1) the magnetic field intensity measured by Neubauer's instrument increased from $\approx 70$ to $\approx 15 \gamma$ within 2 min (it increased from 7.7 to 11.5 $\gamma$ in 64 s) and (2) the plasma speed density and temperature increased between 0122 and 0205 UT (see Figures 2 and 3). The shock normal computed from the magnetic field data using the coplanarity theorem is $\lambda_n = 60°$, $\theta_n = 14°$, which is close to that expected for corotation in a 300-km/s wind, namely, $\lambda_n = 50°$, $\theta_n = 0°$; here $\lambda_n$ is the heliographic longitude, which is
Fig. 3. The corotating stream interface (top) seen by each of the spacecraft. The interface is defined by the abrupt decrease in density and the corresponding increase in temperature. Times have been shifted so that the interfaces are aligned vertically, allowing a comparison of the density, temperature, and magnetic field intensity profiles (bottom).

taken to be zero for a vector pointing radially away from the sun, and \( \theta_n \) is the latitude with respect to the ecliptic plane. The shock speed computed from the observed densities and bulk speeds using the coplanarity normal is 300 km/s, or 540 km/s in the radial direction. This implies that the shock should have arrived at earth 41 hours after it passed Helios 2 (if it moved at constant speed), that is, at hour 19 on November 26. A ssc was reported at 1704 UT on November 26, in good agreement with the prediction. Imp 8 was in the solar wind on November 26, but there are data gaps at the time of the ssc. Nevertheless, the magnetic field intensity nearly doubled at some time in a 2-hour interval centered about the ssc (Figure 3), and the plasma density, temperature, and speed increased at some time in a 5-hour data gap which included the time of the ssc (Figures 2 and 3). Thus the Imp 8 data are consistent with the presence of a shock at earth at 1704 on November 26. Altogether, the data from Helios 2, Imp 8, and the ssc give fairly convincing evidence for a shock driven by a corotating stream, which moved nearly radially from 0.6 to 1 AU and on to Voyager 1, 2 at 1.6 AU. Figure 3 shows that the shock moved away from the interface during the time that it moved from 0.6 to 1.6 AU.

It is customary to refer to a shock ahead of a corotating stream as a corotating shock. This is not appropriate for shock

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<th>TABLE 2. Stream Interface Corotation</th>
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<td>From</td>
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<td>-----------------------</td>
</tr>
<tr>
<td>Helios 1 (Nov. 23, 0245)</td>
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<tr>
<td>Helios 2</td>
</tr>
<tr>
<td>Imp</td>
</tr>
<tr>
<td>Voyager 1</td>
</tr>
</tbody>
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B, however. If shock $B$ were corotating, then it should have been detected at Helios $1 \pm 50$–60 hours before it was observed at Helios $2$ (i.e., late on November 23), because Helios $1$ was at the same radial distance as Helios $2$ and $\approx 35^\circ$ to the east. Although the Helios $1$ observations are nearly complete and continuous, there is no evidence of a shock at Helios $1$ (see Figures 2 and 3). A possible explanation is that the stream which produced the shock was corotating but not stationary. For example, the stream may have been produced by a coronal hole that rotated with the sun but whose physical characteristics changed on a scale of 1 day, producing a time-varying stream profile. Indeed, Figure 2 shows that the speed profile measured by Helios $2$ differs in some details from that measured by Helios $1$, indicating some time variations in this case. Evidence for nonstationary, corotating streams was presented earlier by Burlaga et al. [1978]. (A model of nonstationary flows was presented by Wu et al. [1979], but this does not include the dynamical effects of the sun’s rotation.) Shock $B$ was seen at Helios $2$, Imp $8$, and Voyager $1$, $2$ because those spacecraft were near a radial line; once formed at $\leq 0.6$ AU, the shock persisted and was convected past the other spacecraft. But apparently, conditions were different at the time the stream was at Helios $1$, $35^\circ$ east of Helios $2$, and were not favorable for the production of a shock.

3. Events Associated With a Flare

On November 22, 1977, a $B$ flare at N23, W40 in McMath plage region 15031 was observed in H$\alpha$ starting at 0946 UT and reaching maximum intensity at 1006 UT. Chambon et al. [1978] observed hard X rays and $\gamma$ rays from the flare starting at $\approx 1000$ UT. It produced a SID, a type IV burst (starting at 1002), a type III burst (beginning at 0959 UT), an interplanetary shock wave, and an energetic particle event. Thus the event displayed a wide range of phenomena that one associates with a great flare [Dryer, 1974].

Type III bursts. The type III solar radio burst produced by the flare is the most intense observed to date by Helios $1$ and $2$. Helios $2$ radio observations of the November 22 burst are shown in Figure 6. They are from the University of Minnesota (52, 77, and 203 kHz) and Goddard Space Flight

Fig. 4. Structure of the interface, shown by a plot of high-resolution magnetic field and plasma data (top) and corresponding plasma wave observations. The interface is relatively broad (30 min), its structure does not change appreciably over the 0.2-AU separation between Voyager $1$ and Voyager $2$, and there is no evidence of an instability that might produce waves $\leq f_p^*$ at Voyager $2$.

Fig. 5. Shock $B$, showing the high-resolution magnetic field and plasma data (top panel) and plasma wave observations (bottom panel) near the shock. The flow and field parameters are steady before and after the shock front, allowing accurate determination of its normal and speed. Whistler wave turbulence is observed at $f < f_p^*$ behind the shock; a short burst of broadband turbulence is observed at the shock; and 'ion acoustic' waves are observed at $f_p^* < f < f_p^-$ ahead of the shock.
Center experiments. Electron observations from the Max-Planck-Institut fur Aeronomie experiment are also displayed in Figure 6, showing that electrons in and near the 20- to 65-keV energy range were present, consistent with the idea that low-frequency type III solar radio emission is caused by electrons with energies 10–100 keV [Lin et al., 1973]. Despite the data gap around 1010 UT, it is clear that the radio burst was double peaked at the higher frequencies, possibly due to two separate bursts; however, there was only a single peak at lower frequencies. The first peak reached maximum intensity at 1001 UT for 3 MHz, and the merged peak is observed at 1032 UT for 77 kHz. Much of this delay corresponds to the transit time for the energetic electrons from a heliocentric distance of 0.05 AU (3-MHz level) out to 0.8 AU (77-kHz level), indicating an outward speed greater than 0.2c for the exciter. A few minutes of the delay arise from the difference in propagation time of the electromagnetic waves from the source levels to Helios 2, located at 0.6 AU.

Flux densities observed for this burst by Helios 2 reached maximum values exceeding $10^{-15}$ W m$^{-2}$ Hz$^{-1}$ for frequencies from 77 to 255 kHz; they decreased to approximately $10^{-16}$ W m$^{-2}$ Hz$^{-1}$ at 3 MHz, the highest Helios observing frequency. The 52-kHz channel, which shows strong electrostatic noise from 1025 to 1050 UT, is at the peak of the electrostatic noise spectrum, and it is within 1–2 kHz of the local plasma frequency determined from the measured density. Similar bursts were reported by Garnett et al. [1978] and Garnett and Anderson [1977]. The electrostatic bursts might be short in comparison to the sampling time of the tuned receiver. The bandwidth of the receiver is about 5 kHz, and its rise time therefore is about 0.2 ms, which is instantaneous in comparison to the detector integration time of 50 ms. As a consequence, for signals whose duration is more than 0.2 ms, the measurement gives the input voltage averaged over 50 ms.

The 77-kHz channel is the lowest frequency which did not show electrostatic noise. Burst radio emission has been reported to be generated at twice the local plasma frequency [Alvarez et al., 1972]. Consequently, the 77-kHz electromagnetic waves detected at 0.60 AU by Helios 2 were propagating backward toward the sun from a source level near 0.8 AU, where the plasma frequency is half of 77 kHz.

This burst and the associated electron beam were also observed by the Voyager 1 and 2 planetary radio astronomy experiment and low-energy particle experiment, respectively. The burst arrival directions, found by the spinning Helios 1 and 2 antennas, together with the Helios and Voyager electron data, show that the exciter extends over a wide (＞75°) range of solar longitudes. Analysis of the relative intensities and positions observed by Helios 1 and 2 also indicates that the centroid of the burst passed between these two spacecraft. Assuming a source longitude of 40°W and a spiral field configuration, a best fit to the intensity versus frequency data obtained by Helios 1 and 2 is obtained for a solar wind speed of 300 km/s. This is consistent with the speeds measured by the Helios plasma instruments, which were near 300 km/s for several days.

Interplanetary shocks and flows. The interplanetary shock wave produced by the flare was observed directly by Helios 1 and 2, Imp 8, and Voyager 1 and 2; it was also observed indirectly as a scs at the earth (see Table 4 and Figure 7). The shock might have been driven by ejecta, as is suggested by the sketch in Figure 8, but the ejecta were not actually observed, because no spacecraft was suitably positioned.

If one tries to determine the motion of the shock by using a radial distance versus time plot (Figure 9) and the customary assumption of spherical symmetry, one encounters difficulties that would have been overlooked if there were fewer spacecraft. One difficulty is that the speed determined from the time delay between Imp 8 and Voyager 1 is 418 km/s, whereas the speed determined from the time delay between Imp 8 and Voyager 2 is 568 ± 20 km/s (the uncertainty is due to a data gap at Voyager 2 between 0600 and 0900 UT). This discrepancy is large, considering that Voyager 1 and Voyager 2 were separated by only 0.005 AU in the radial direction and by 0.2 AU in the transverse direction.

A second and more extreme example of the inadequacy of the assumption of spherical symmetry for computing shock speeds is the speed determined from the time delay between Voyager 1 and Voyager 2: 14 ± 2 km/s! This is obviously
TABLE 4. Shock A

<table>
<thead>
<tr>
<th>Shock</th>
<th>Helios 2</th>
<th>Helios 2</th>
<th>Helios 1</th>
<th>Imp 8</th>
<th>Voyager 2</th>
<th>Voyager 1</th>
</tr>
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<tr>
<td>Date</td>
<td>A₁</td>
<td>A₂</td>
<td>A₃</td>
<td>A₄</td>
<td>A₅</td>
<td>A₆</td>
</tr>
<tr>
<td>UT</td>
<td>Nov. 23</td>
<td>Nov. 24</td>
<td>Nov. 25</td>
<td>Nov. 25</td>
<td>Nov. 27</td>
<td>Nov. 27</td>
</tr>
<tr>
<td>Time, UT</td>
<td>1610</td>
<td>0611</td>
<td>2228</td>
<td>1213</td>
<td>0730 ± 0130</td>
<td>2226</td>
</tr>
<tr>
<td>(r_1, 10^8 ) km</td>
<td>0.916</td>
<td>0.927</td>
<td>1.018</td>
<td>1.476</td>
<td>2.361</td>
<td>2.533</td>
</tr>
<tr>
<td>Normal (\lambda_n), deg</td>
<td>16</td>
<td>-15</td>
<td>4</td>
<td>-34</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Normal (\theta_n) or (\delta_n), deg</td>
<td>(\theta_n = -14)</td>
<td>(\theta_n = -48)</td>
<td>(\theta_n = 25)</td>
<td>(\delta_n = -10)</td>
<td></td>
<td></td>
</tr>
<tr>
<td>(V_n), km/s</td>
<td>330</td>
<td>304</td>
<td>352</td>
<td>302</td>
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<tr>
<td>(V_n) km/s</td>
<td>353</td>
<td>467</td>
<td>390</td>
<td>(-418)</td>
<td>369</td>
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wrong, and it is far from the speed determined from the analysis of the shock data at Voyager 1 (Figure 10), namely, 302 km/s. The shock normal and speed computed from the Voyager 1 data by using the method of Lepping and Argentiero [1971] were \((\lambda_n = -34^\circ, \theta_n = -10^\circ)\) and \(V_n = 302 \) km/s, respectively (see Table 4). Using these numbers, assuming that the shock was plane between Voyager 1 and Voyager 2, and considering the inertial solar ecliptic positions of Voyager 1 \((r_1 = (2.280, -0.274, 0.115) \times 10^8 \) km) and Voyager 2 \((r_2 = (2.285, -0.533, 0.267) \times 10^8 \) km), one finds that the predicted time delay between Voyager 1 and Voyager 2 is 11 hours 13 min, which is reasonably close to the observed delay, \((15 \pm 1.5)\) hours. (The ±1.5-hour uncertainty is due to a data gap.) The small discrepancy may be attributed to uncertainties in the shock normal and to curvature of the shock surface. By contrast, the time delay predicted by using the assumption of spherical symmetry is only 36 min. We conclude that the use of time delays and assumption of spherical symmetry do not always give accurate shock speeds, whereas the use of local jump conditions and observations did give reasonably accurate estimates of the shock speed and direction in this case. The observed orientation of the shock is consistent with that

Fig. 7. Shocks A₁, A₂, and A₃ and the stream interface. At 0600 UT, November 24, 1977, the interface had passed Helios 1 but had not reached Helios 2. One shock (A₁) was approaching Helios 1 and arrived at Helios 1 late on November 24. Two shocks were observed by Helios 2. One (A₂) arrived at Helios 2 at 0611 UT on November 24, and another was a short distance ahead of it. A₁ and A₂ coalesced into one shock (A₄) as they moved from Helios 1 to earth, where A₄ was detected by Imp 8.
Fig. 8. A sketch, approximately to scale, showing the position of shock B, the stream interface, and shock A₂ at 0600 UT, when A₂ was approaching Helios 2. The positions of the spacecraft and the flare site at that time are also shown. The hypothetical ejecta were not observed. The flare produced energetic protons which escaped freely through the stream. Shock A₂ accelerated particles locally and perhaps trapped some of the flare particles, producing a local maximum in counting rates at the shock observed by Helios 2. These shock-accelerated particles did not penetrate the stream interface and were not observed by Helios 1.

expected for a shock with a radius of curvature less than 1.6 AU, originating at the flare site.

Helios 2 observed two shocks (A₁ and A₂, at 1610 UT on November 23 and at 0611 UT on November 24, respectively (see Figures 7 and 11)). However, Imp 8, which was at nearly the same latitude and longitude which was only 0.36 AU away from Helios 2, observed only one shock (A₄ at 1213 UT on November 25 (see Figures 7 and 10)). We cannot unambiguously determine why two shocks passed Helios 2 (several origins can be imagined), but we can suggest why only one shock was subsequently observed at Imp 8 and at Voyager 1 and 2. The radial speed of shock A₁, determined from the local plasma and magnetic field observations of the shock by method MD1 of Abraham-Shrauner and Yu (1976), was 353 km/s. The corresponding speed of A₂ was 467 km/s. Thus although A₂ followed A₁ (i.e., it was closer to the sun (see Figure 7)), it was moving faster than A₁. Consequently, A₂ should have overtaken A₁ at some point; assuming constant speeds, this point was at 1.08 AU on the Helios 2-sun line. If the computed shock normals (Table 4 and Figure 7) are even approximately correct, the shocks should have interacted along the earth-sun line before they reached Imp 8 near the earth. The observation of only one shock at Imp 8 suggests that when the shocks interacted, they coalesced. This is in agreement with gas dynamic theory, where the overtaking of one shock (A₁) by a following one (A₂) leads to a coalesced shock moving forward and a reverse rarefaction fan, which because of its spreading is difficult to observe. (In MHD the interaction leads to seven distinct MHD structures, the most prominent ones of which are a forward fast shock and a reverse fast rarefaction wave.) The resultant shock propagated to Voyager 1, which was close to the earth-sun line. Its radial speed at Voyager 1, which was close to the earth-sun line. Its radial speed at Voyager 1, determined from the shock observations by using the method of Lepping and Argentiero (1971), was 369 km/s, which is in reasonable agreement (considering typical normal errors) with the speed determined from the time delay between Imp 8 and Voyager 1, namely, 427 km/s. Evidence for shock coalescence in Pioneer data has been reported by Smith et al. (1977). A second alternative would be a sufficient weakening of one shock before it interacted with the other one. (Note that the very weak shock would still have to interact with the second shock.) This possibility is ruled out by the following two arguments: A₁ cannot be the weakened shock, since it fits very nicely into the propagation diagram (Figure 9) in contrast to shock A₁. We rule out a large weakening of shock A₁, since it is followed by a long-lasting region of increased momentum and energy flux as shown in Figure 7.

A remaining aspect which requires clarification is the observation of one shock only at Helios 1. A possible explanation for this observation may be the presence of the stream interface and an interaction region between Helios 1 and Helios 2 (Burlaga and Scudder, 1975). If we approximate the interface as a tangential discontinuity, its interaction with A₁ may lead to the latter's disappearance [e.g., Neubauer, 1976].

4. The December Shock

During the Helios-Voyager-Imp workshop it was noted that a shock was observed by Helios 1 and 2 on December 1 and by Voyager 1 and 2 on December 2, and it was decided to include this event in the joint study. The interplanetary data are nearly complete for Helios 1, as is shown in Figure 12. However, the solar data do not show any large flare which might have produced the shock. One candidate is an SN flare at S24, E85 which began in H₃ on 0338 UT on November 30 and reached a maximum at 0350 UT. This small flare was associated with an X-ray burst (starting at 0330 UT, with a maximum at 0348 UT) and a SID (starting at 0334 UT, with a maximum at 0349 UT). This implies deceleration of the shock within 0.6 AU [see Gosling et al., 1968]. In view of the uncertainty concerning the source of the shock our discussion emphasizes the interplanetary observations.

Postshock conditions. The density and temperature profiles in Figure 12 suggest that the shock observed by Helios 1 was followed by ejecta in which the density and temperature were low. There is also evidence for enhanced magnetic field intensities in the ejecta. Helios 2 may also have observed the

Fig. 9. Propagation of shock A. The radial position of the shock is shown at the times that the shocks arrived at Helios 1 (H₁), Helios 2 (H₂), Imp 8, Voyager 1 (V₁), and Voyager 2 (V₂). The two shocks observed by H₂ coalesced into the one shock observed by Imp 8. Departures from spherical symmetry of shock A are indicated by the scatter of the points about a straight line.
Fig. 10. High-resolution magnetic field and plasma data showing that $A_4$ (at Imp 8) and $A_6$ (at Voyager 1) are shocks. A narrow, broadband burst of electrostatic noise was observed at the time of the shock by both spacecraft. 'Ion acoustic' waves between $f_-$ and $f_+$ were observed upstream by both spacecraft. Voyager 1 also observed whistler mode turbulence at $f < f_+$ behind the shock.

The shock was also detected by Voyager 1 and 2 (see Figures 13 and 14), but they did not encounter ejecta like that seen by Helios 2. Voyager 1 observed a small depression in density on December 4 and a small increase in B on December 4-5; this was probably a local phenomenon, since it was not seen by Voyager 2, which was near by. This signature is clearly different from that seen by Helios 1, so we do not identify it as ejecta.) Thus the evidence is that the shock had a wide longitudinal extent ($\geq 40^\circ$ (see Figure 15)) and was driven by ejecta less broad, originating east of the Voyager-sun line.

Note that Voyager 2 observed a nearly monotonic decrease in speed, density, temperature, and magnetic field strength behind the shock (Figure 13). Many authors have interpreted such a signature as evidence for a blast wave, generally on the basis of observations from just one satellite [see Hundhausen, 1972]. However, the observation of ejecta at Helios 1 indicates that this was probably not a blast wave; it was a driven shock. Voyager 2 saw the shock, but it did not encounter the ejecta owing to its more limited longitudinal extent. This shows that the signature of the postshock flow is not sufficient to identify the type of a shock wave. This point was made previously by Ogilvie and Burlaga [1974] and by Rosenau and Frankenthal [1978]; it has recently been demonstrated very clearly by Acuna et al. [1980]. The concept of a broad shock driven by narrow ejecta is not new, although it is often forgotten or ignored. It dates back at least to Gold [1959].

Shock motion. Figure 16 gives a plot of radial distance versus time, showing the shock positions and times determined from the observations of Helios 1, 2 and Voyager 1, 2 and from a sudden commencement at earth. The points lie very close to a straight line with a slope corresponding to a speed of 555 km/s. Considering that Helios 1 was 19° east of the Voyager 2-sun line and that earth was 17° west of that line the straight line in Figure 16 suggests a nearly spherical
Fig. 11. High-resolution magnetic field and plasma data showing that $A_1$ and $A_2$ are shocks (or steep compressive waves). Electrostatic plasma wave data from Helios 2 show that the shock was imbedded in a broad region of Doppler-shifted 'ion acoustic' waves. A narrow spike was observed at 562 and 311 Hz at the time of the shock.

A shock front moving at a constant speed between 0.6 and 1.6 AU. Similar results for the August 1972 events were reported by Smith et al. [1977] and Dryer et al. [1976]. However, examination of the local shock speeds and normals reveals a more complicated picture. Since Voyager 2 and Helios 2 were nearly radially aligned and since Figure 16 suggests a spherical shock, one expects that Voyager 2 and Helios 2 should have observed essentially the same shock speed and direction, the radial component of velocity being close to 555 km/s. The local jump conditions give rather different results (Table 5): (1) the local speeds were substantially less than the speed determined from the average speed determined from the time delay; and (2) the shock normal at Helios 2 ($\lambda_n = -3^\circ$, $\theta_n = 17^\circ$) was very different from that at Voyager 2 ($\lambda_n = 38^\circ$, $\theta_n = -6^\circ$). These differences are too large to be attributed to uncertainties in the computation of the local shock speed and direction. The field and plasma parameters were relatively steady before and after the shock, the field direction change was relatively large (18$^\circ$ at Helios 2), and we used both magnetic field and plasma observations, so we expect the uncertainty in speed to be $\approx 20$ km/s and the uncertainty in direction to be $\approx 10^\circ$ [Abraham-Shrauner and Yun, 1976; Lepping and Arge ntiero, 1971]. Thus the observations suggest that locally the shock surface may have been distorted such that the normal was not radial, although the normal may have been radial on average. Likewise, locally the shock may have been accelerated or decelerated, giving local speeds higher than average in one place, lower than average in a second place, and near average in a third place [Heinemann and Siscoe, 1974; Burlaga and Scudder, 1975]. For example, the radial component of the local velocity at Voyager 2 (530 km/s) is consistent with the average speed determined from time delay (555 km/s) within the experimental uncertainties, but the radial component of the local velocity at Helios 2 (460 km/s) is substantially less than the average value. Since Helios 2 and Voyager 2 were nearly radially aligned, this suggests that the radial component of the shock velocity may have fluctuated as much as $\approx \pm 100$ km/s and its direction may have fluctuated as much as $\pm 40^\circ$ as it moved between 0.6 and 1.6 AU. The alternative is to postulate very large azimuthal variations.

5. Plasma Waves at Shocks

Helios 1, 2 and Voyager 1, 2 carried plasma wave instruments (see Garnett and Anderson [1977] and Scarf and Garnett [1977], respectively, for a discussion of the instruments), which provided an extensive set of observations of waves near the interplanetary shocks discussed above. These observations were used as a means of searching for and confirming the identity of the shocks. More important, however, they provide an exceptionally large and complete record which forms a basis for a comparative study of waves at interplanetary shocks. Only a few papers discussing plasma wave electric fields at interplanetary shocks have been published [Scarf, 1978; Scarf et al., 1979; Garnett et al., 1979a, b].

Fig. 12. High-resolution plasma and magnetic field data showing shock C and a boundary behind it, which might be the boundary of the ejecta (contact surface). Note the depression in magnetic field intensity at the boundary. Electrostatic plasma waves are observed between $f_p^+$ and $f_p^-$ at the shock, but no significant waves are observed at the contact surface.
we shall present only a qualitative discussion stressing the remarkable variety of signatures. A more comprehensive physical discussion is deferred to another paper.

The wave data are given together with the plasma and magnetic field observations of the shocks in Figures 5, 10, 11, 12, and 14. The electric field intensity is plotted versus time for each of several frequency channels on a logarithmic scale with a range of 100 dB for each channel. The electric field strength ranges from about 1 µV m⁻¹ at the bottom of the scale to 100 mV m⁻¹ at the top of the scale. The solid lines represent peak electric field amplitudes, and solid black areas (or vertical solid lines in some cases) represent the average electric field amplitude.

Let us consider the individual shock observations in the order in which shocks were introduced above, beginning with shock B. This shock had not developed at the position of Helios 1, but it was observed at both Voyager 1 and Voyager 2 (Figure 5), which were at essentially the same radial distance (1.6 AU) and separated by ~0.2 AU. The Voyager 1 plasma wave observations show at least three different types of emissions: (1) turbulence extending downstream of the shock at frequencies of ≲ f₁⁺, identified as whistler mode turbulence; (2) waves extending upstream at frequencies from about 1.0 to 5.62 kHz, tentatively identified as ion acoustic waves; and (3) a short, well-defined broadband burst at the shock at frequencies from 10 Hz to 5.62 kHz. These types of emissions have been discussed by Scarf et al. [1970], Garnett and Frank [1978], and Garnett et al. [1979a]. Voyager 2 also observed the whistler mode turbulence extending downstream from the shock, and it observed a peak corresponding to the broadband emissions at the shock. There are no Voyager 2 data above 1 kHz, probably because of a failure in the spacecraft data system which reduced the sensitivity of these channels.

Plasma waves at shock A were observed by Helios 2 (Figure 11) and by Imp 8 and Voyager 1 (Figure 10). Whistler waves were not observed downstream of the shock at Helios 2 and Imp 8, but they were observed downstream of the shock at Voyager 1. The shock at Helios 2 is almost totally obscured by a broad region of ion acoustic wave turbulence from about 562 Hz to 10 kHz; these waves are not necessarily all associated with the shock [Garnett and Frank, 1978]. Imp 8 and Voyager 1 (Figure 10) observed ion acoustic waves upstream of the shock between f₁⁺ and fₚ⁻. Helios 2 (Figure 11) observed a sharp burst of noise in the 311- and 562-Hz channels coincident with the passage of shock; Imp 8 found some evidence of a corresponding noise burst below f₁⁺, and Voyager 1 observed a noise burst at the shock in the range 31 Hz to ≈1.78 kHz.
Shock C was observed by Helios 1 and 2 and by Voyager 2 (Figures 12–14). None of the spacecraft observed intense whistler mode turbulence behind the shock. Helios 1 and Helios 2 observed an enhancement in electric field intensity in the range 562 Hz to 10 kHz with a large peak to average ratio, probably due to Doppler-shifted ion acoustic waves [Gurnett and Frank, 1978]. The waves extended both upstream and downstream at Helios 1 but only downstream at Helios 2. Voyager saw only weak emission of such waves, downstream of the shock. A sharp, intense (1–5 mV m⁻¹) broadband burst of electric field turbulence was observed at Helios 2, but it was absent at Voyager 2 and missing or obscured by the ion acoustic waves at Helios 1.

We conclude that at least three types of emissions (in various combinations) may be observed at an interplanetary shock, namely, downstream ‘whistler mode turbulence,’ upstream ‘ion acoustic’ waves, and a brief broadband noise burst coincident with the shock. In some cases, only one or two of these are observed. In addition, the shock may be embedded in a broad region of ion acoustic waves not necessarily caused by the shock. The combination of wave types and the characteristics of each wave mode seen at one spacecraft may be very different from those observed by another spacecraft nearby. Apparently, the plasma waves at a shock depend strongly on the local characteristics of the medium. This is not surprising, since it has been observed in the case of the earth’s bow shock [Greenslade et al., 1973]. However, the basic types of emissions are the same at 0.6 AU as they are at 1.6 AU.

6. ENERGETIC PROTONS

In the interval November 22 to December 6, 1977, Helios and Voyager instruments observed energetic protons (±50–200 keV) produced by at least two mechanisms: local shock acceleration and acceleration in a flare. It is convenient to begin by discussing the former, since shock-accelerated particles are less complicated by propagation effects.

Shock acceleration. Protons accelerated by a shock are seen most clearly in the case of shock C, which was relatively isolated and uncomplicated, as was discussed in section 4. Recall that Voyager 1 and 2 observed a shock behind which the flow parameters and magnetic field intensity dropped gradually to the preshock values; there was no evidence of ejecta like those observed by Helios 1. Figure 17 shows enhancements in the counting rate of protons at Voyager 1 and 2 in the energy range ±50 to ±138 keV; the maximum intensity occurred at or just behind the shock. At Voyager 2 the peak counting rate was ≈100 times the ambient value, and at Voy-
ager 1 the enhancement was somewhat smaller. The enhancement began $\approx 15$ hours ahead of the shock at both Voyager 1 and Voyager 2. It persisted for $\approx 32$ hours behind the shock at Voyager 1 and $\approx 28$ hours behind the shock at Voyager 2. There were small differences in the shapes of the profiles which might be due to differences in the local magnetic field configurations. Basically, however, the proton enhancement at Voyager 1 was similar to that at Voyager 2. This may be due to the simple geometry of the shock near Voyager 1 and 2 and to their relatively small separation (0.2 AU).

The situation at Helios 1 and 2 was quite different. Both spacecraft observed an enhancement in counting rate of protons (Figure 17). The maximum enhancement at Helios 2 was only $\approx 20$ times the background counting rate, and it occurred at the shock. Two maxima were observed by Helios 1, and the shock occurred between them. A compression wave was observed at the time of the second maximum (Figure 12), but the time resolution was not adequate to determine whether or not it was a shock. The counting rate dropped abruptly approximately 6 hours after the shock at both Helios 1 and Helios 2, in contrast to the more gradual, longer-lasting decline at Voyager 1 and 2. This might be due, at least in part, to the presence of ejecta at Helios 1 and at Helios 2, which were not observed by Voyager 1 and 2. (There is no accepted signature for the boundary of ejecta, and we cannot be certain that we have identified one. The vertical line behind the shock in Figure 17 corresponds to an abrupt decrease in density observed behind the shocks in Figures 12 and 13.) The enhancement began $\approx 6$ hours ahead of the shock at Helios 2 and a few hours ahead of the shock at Helios 1; the slight difference could be due to different acceleration efficiencies of the two shocks and/or to different upstream magnetic field conditions which gave connection to the shocks at slightly different times. There is a curious enhancement at Helios 1, occurring several hours ahead of the shock-associated enhancement but closely resembling it. One can imagine that this was due to a magnetic field geometry which provided a good connection between the observer and the shock for several hours before the shock arrived.

The differences between the enhancements at Helios 1 and Helios 2 and the differences between the enhancements at Voyager 1 and Voyager 2 indicate that local conditions do influence the intensity profile somewhat. Note, however, that the Voyager 1, 2 profiles have a greater maximum enhancement and a greater upstream extent than the Helios 1, 2 profiles. One possible reason for this (but not the only one) is that Voyager 1, 2 were farther from the sun than Helios 1, 2, so that the shock at Voyager 1, 2 had been accelerating particles for a longer time and perhaps accelerated and accumulated more particles than it had when it was at the positions of Helios 1 and 2.

Flare-accelerated particles. Now let us discuss the low-energy ($\approx 25$–$200$ keV) protons ejected by the flare of November 22, 1977 (see section 3 for a discussion of the flare characteristics and the corresponding interplanetary flows). Helios 1 and 2 observed very different intensity-time profiles during the decay in intensity (Figure 18), even though they were at nearly the same radial distance and were separated in longitude by only $32^\circ$ (see Figures 1 and 8). At Helios 1 the intensity decreased smoothly and monotonically for at least 3 days (Figure 18). The corotating stream discussed in section 2 was east of Helios 1 at the beginning of the event, and the interface

<table>
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<th>Helios 2</th>
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<th>Imp 7</th>
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<td>$V$, km/s</td>
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<td>501</td>
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Fig. 17. Counting rates of energetic protons near shock C, observed by Helios 1, 2 and Voyager 1, 2. The broad, intense fluxes of particles at Voyagers 1 and 2 closely resemble one another, but they differ appreciably from the narrower, less intense fluxes observed by Helios 1 and 2. The Helios 2 profile differs appreciably from that of Helios 1. The abrupt decrease in counting rates behind the shock observed by Helios 1 and Helios 2 may be due to a flow boundary (e.g., a piston) behind the shock.
of the decay seen by Helios 2 was very similar to that observed by Helios 1 (Figure 18), the flux decreasing monotonically for at least 12 hours. As shock A₁ (produced by the flare) approached Helios 2, the counting rate of energetic protons began to increase, reaching a maximum at the time that shock A₂ reached Helios 2. The maximum flux was $4 \times 10^9$ ions cm$^{-2}$ s$^{-1}$ sr$^{-1}$ MeV$^{-1}$. The maximum counting rate was $\approx 25$ times that measured by Helios 1 at the same time, that is, comparable to the increase which Helios 2 observed at shock C, as described above. This increase may be due to (1) particles accelerated by the shock, (2) flare particles trapped behind the shock, and/or (3) energetic storm particles. Following the shock the counting rate again decreased until the interaction region of the corotating stream arrived at Helios 2, at which time there was a slight increase in the counting rate, perhaps due to compression in the interaction region. When the interface arrived, the counting rate at Helios 2 dropped rapidly (exponentially with a time scale of 3 hours) to approximately the same level that Helios 1 recorded. Apparently, particles accelerated by shock A could not penetrate the stream interface, and many were trapped in a region bounded by the interface on one side and the shock on another side. The ejecta from the flare might have provided the third boundary. The scenario that has been described is represented schematically in Figure 8.

Voyager 1 and 2 observed intensity-time profiles of protons in the energy range $\approx 50$–138 keV (Figure 19) which resemble the profile recorded by Helios 2. During the early stage of the decay both spacecraft observed a monotonic decrease in counting rate lasting $\approx 16$ hours. (The initial increase in counting rate and the first hour or two of the decay include an uncertain contribution to energetic, omnidirectional particles.) The counting rate then increased gradually during the next 8 days, reaching a maximum counting rate at the time that shock A arrived. (Recall that there was a data gap at Voyager 2 between 0600 and 0900 UT, so the shock was not observed directly.) This gradual increase lasted too long to be due to particles accelerated by the shock alone. Probably, energetic storm particles were present. The rapid increase several hours ahead of the shock at Voyager 1 and 2, however, is probably a contribution due to shock acceleration. The enhancement is relatively small, no more than about 16 times the ambient...

Fig. 19. Counting rates of energetic protons observed by Voyager 1 and 2. Speed profiles are shown to indicate the flow conditions. Both Voyager 1 and Voyager 2 observed flare particles on November 22. The broad increase between November 27 and 29 may be due to energetic storm particles. Locally accelerated particles are observed at the transient shock A, but there is no significant increase at the stationary shock B.
value. It did not extend above 0.5 MeV for protons. No modulation of electrons in the range 0.03–1.5 MeV was observed. At the time of the shock, Voyager 1 observed a strong anisotropy (3.5:1), the particles flowing away from the sun. Shortly after the shock passed, the anisotropy direction reversed, and particles were observed to be streaming toward the sun, consistent with the hypothesis that most of the particles observed near the shock were accelerated by the shock. Following the shock the counting rate decreased, rapidly at first and then more slowly. Shock B (see section 2) arrived at Voyager 1 and 2 during the decline in intensity, on November 29, and the corotating stream interface arrived several hours later. A very small increase in the counting rate of low-energy protons was observed by Voyager 2, and an even smaller increase by Voyager 1, but these were insignificant in comparison to the other shock-associated enhancements described above. A small increase in counting rate was observed in the interaction region ahead of the interface (Figure 19), analogous to that observed on November 25 by Helios 2 when it encountered the interaction region (Figure 18).

7. Summary

We have presented a wealth of data obtained at $\approx 0.6$, 1, and 1.6 AU, describing the evolution and interactions of particles, flows, and fields in the period November 22 to December 6, 1977. Some of the principal results of our analysis of these data are as follows:

1. A small, corotating stream, originating in a coronal hole, was observed to disappear as it moved from 0.7 to 1.6 AU. A forward shock (shock B) was produced by the stream and observed by Helios 2 (0.6 AU), Imp 8 and earth (1 AU), and Voyager 1, 2, which were nearly radially aligned; however, the shock was not corotating because it was not seen at Helios 1, 35° east of Helios 2. Apparently, the shock was corotating but nonstationary. The stream interface corotated from 0.7 to 1.6 AU and persisted even though the stream had dissipated; it was stable, and its structure remained essentially the same at all positions.

2. An exceptionally intense type III burst, produced by the November 22, 1977, flare, was observed by Helios 1 and 2. The electron beam which caused it and plasma oscillations excited by the beam were observed at 0.6 AU.

3. The shock produced by the flare of November 22 (shock A) was nonspherical, pointing 34° to the east and 10° south of the radial direction at 1.6 AU. It interacted with another shock beyond 0.6 AU, and they coalesced, forming a single shock that was observed at 1 and 1.6 AU.

4. A shock of uncertain origin (shock C) was observed by five spacecraft at radial distances from the sun ranging from 0.6 to 1.6 AU and with longitudinal separations of up to 36°. The radial distances versus time diagram suggested a spherical shock moving at a constant speed, but analysis of data at the shocks showed local fluctuations of up to 100 km/s in speed and 40° in direction.

5. One or more of three types of electrostatic waves were observed at interplanetary shocks: upstream waves with $f < f_p$, downstream waves with $f < f_p$, and broadband noise at the shock. These three types of emission were observed at 1.6 AU as well as 0.6 AU. The specific pattern varied greatly among the shocks observed, even for the same shock observed at closely separated ($\approx 0.2$ AU) spacecraft, indicating a strong dependence on local shock and solar wind parameters.

6. Energetic protons ($\approx 50$–200 keV) were observed to be accelerated at shocks. The maximum and half widths of the flux profiles at a shock differed by approximately a factor of 2 over distances of a few tenths of an astronomical unit, indicating a dependence on local conditions. The data suggest a tendency for the fluxes to become broader and more intense with increasing distance from the sun.

7. Energetic protons ($\approx 50$ keV) from the November 22, 1977, flare were observed. Helios 1 observed that their intensity decayed monotonically in the corotating stream, with little change across the stream interface. Helios 2, 30° to the west of the interface, observed a very different profile, with a second increase to a maximum at the time that the shock produced by the flare arrived. These ‘extra’ particles apparently did not penetrate the interface, for the intensity at Helios 2 dropped abruptly to the intensity observed at Helios 1 when the interface corotated past Helios 2.

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